angle not too close to 45° is ideal for obtaining information about perfection by the polarization method. This is so because for a strong reflection, α_{PA} is not reduced much below the value of α_{PN} by absorption and the full range from $|\cos 2\theta|$ to $\cos^2 2\theta$ is available.

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The Crystal Structures of some Anthracene Derivatives VI. 9-Anthraldehyde

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Crystals of 9-anthraldehyde are orthorhombic, space group $P2_12_12_1$, with four molecules in the unit cell. The structure has been determined from the projection along the short a-axis, and refined by three-dimensional differential syntheses. The anthracene skeleton is completely planar, but not coplanar with the plane of the aldehyde group, the angle between the planes being 27° .

Introduction

The physical properties of nitromesitylene (I), such as Raman and infrared frequencies, dipole moment and ultraviolet absorption spectrum, indicate that resonance interaction between the nitro group and the aromatic ring is markedly reduced in comparison with nitrobenzene, suggesting that the steric effect of the ortho methyl groups prevents the attainment of a completely coplanar configuration (see, for example, Wheland, 1955). Recently this deviation from planarity has been confirmed and in addition accurately measured by a detailed X-ray analysis of crystals of nitromesitylene (Trotter, 1959 a, b); the nitro group is rotated about the C-N bond 66° out of the plane of the aromatic ring, and this measured angle can be correlated with the characteristic NQ_2 vibration frequencies (Trotter, $1959f$). The X-ray investigations have also been extended to 9-nitroanthracene (II) and 9:10-dinitroanthracene (Trotter, 1958, 1959c, d), where the environments of the nitro groups are very similar to those in nitromesitylene, and the measured deviations from coplanarity have been correlated with the

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infrared frequencies and ultraviolet spectra (Trotter, $1959e, f$).

Spectroscopic and dipole moment evidence indicates similar deviations from coplanarity in acetyl derivatives, and although the deviations in this type of molecule have not been measured by X-ray methods, there is no doubt that in a molecule such as 2:4:6 trimethylacetophenone (III) (acetylmesitylene) the $-CO.CH₃$ group is rotated at least 66[°] out of the aromatic plane. In the case of aldehydes, however, the evidence is much more confused. Raman spectra and dipole moment data have been interpreted as indicating no deviation from coplanarity (Saunders, Murray & Cleveland, 1941; Kadesch & Weller, 1941), but consideration of van der Waals separations would suggest that the -CHO group must be twisted markedly out of the aromatic plane in such molecules as mesitylene aldehyde (IV) and 9-anthraldehyde (V) (Jones, 1945). It might be noted though that in these molecules the strain in a planar model could be relieved somewhat by a small in-plane displacement of the C(aromatic)-C(aldehyde) bond, as well as by an out-of-plane displacement of the oxygen atom. More recent investigations of the ultraviolet spectra of molecules of this type have suggested angles of about 20° between the planes of the aldehyde groups and the aromatic rings (Braude & Sondheimer, 1955), and these authors also show that a 20° twist is not incompatible with the dipole moment data.

A direct investigation of the structure and configuration of a molecule of this type would be useful; a suitable molecule for X-ray examination is 9 anthraldehyde, and its crystal structure has now been determined by X-ray diffraction methods.

Experimental

Crystals of 9-anthraldehyde are yellow in colour, and consist of needles elongated along the a -axis. The density was determined by flotation in aqueous potassium iodide solution. The unit-cell dimensions and space group were determined from rotation and oscillation photographs, and zero and first layer Weissenberg films of crystals rotating about the a and b axes, and from *hO1* and *hkO* precession films.

Crystal data

9-Anthraldehyde, C₁₅H₁₀O Molecular weight = 206.2 ; melting point = $104 °C$. Orthorhombic, $a=4.29 \pm 0.01$, $b=14.11 \pm 0.04$, $c = 17.18 \pm 0.04$ Å. Volume of the unit cell $= U = 1039.9$ Å³. D_m (measured density) = 1.304 g.cm.⁻³, $Z=4$, D_x (calculated density) = 1.309 g.cm.⁻³. Absorption coefficients for X-rays, $\lambda = 1.542$ Å

 μ =7.54 cm.⁻¹; λ =0.7107 Å, μ =0.95 cm.⁻¹.

Total number of electrons in the unit cell = $F(000)$ = 432. Absent spectra: $h00$ when h is odd, $0k0$ when k is odd, 00l when l is odd. Space group is $P2_12_12_1-P_2^4$.

The intensities of the *hkl* reflexions were recorded on Weissenberg photographs for crystals rotating about the a and b axes, using $Cu K_{\alpha}$ radiation, the multiple-film technique being used to correlate strong and weak reflexions. For the crystal rotating about the a-axis layers with $h = 0, 1, 2, 3$ were recorded, and for the crystal rotating about *b* layers with $k=0, 1, 2, 3$, the equi-inclination method being used for all nonequatorial layers. In the *Okl* and *Ikl* zones only about 45% of the reflexions within the Cu K_{α} sphere were observed; the percentage of observed reflexions became smaller as h increased, and only a small number of reflexions with $h=3$ and 4 were observed (only the region of the $4kl$ zone with $k < 4$ has been explored. but all the observed reflexions in this zone are very weak). All the intensities were estimated visually, the range being about 2500 to 1. The values of the structure amplitudes were derived by the usual formulae for a mosaic crystal, Lorentz and polarization factors being applied together with the rotation factors appropriate to equi-inclination Weissenberg photographs (Tunell, 1939). The absolute scale was established later by correlation with the calculated structure factors. Small crystals were used and no absorption corrections were applied.

Structure analysis

[100] *projection*

In deriving an approximate trial structure it was assumed initially that the molecule was completely planar with C-C bond lengths 1.40 Å, C-O bond length $1-20$ A, and all valency angles 120° . The molecular orientation in the a-axis projection was determined from an examination of the $0kl$ weighted reciprocal lattice, and the position of the molecular centre was then found by the usual methods (Taylor, 1954; Trotter, 1959g). The trial structure deduced was checked by interpreting the peaks corresponding to multiple intramolecular and intermolecular vectors on the *Okl* Patterson map. Taking the origin of the projection on a two-fold screw axis (so that the B parts of the structure factors are zero) structure factors were calculated for about 40 of the stronger reflexions, and a Fourier series then computed using measured structure amplitudes and calculated signs. On the resulting electron-density map all the carbon atoms and the oxygen atom were well resolved; the map indicated clearly that the oxygen atom deviates from its idealized position in the trial structure, but from this projection alone it could not of course be deduced whether this involves an in-plane or out-of-plane displacement.

New atomic centres were chosen and structure factors calculated for all the 0kl reflexions, using McWeeny's curves for carbon and oxygen, corrected for thermal vibration as usual, taking $B_c = 5.4 ~ \AA$ ² and

Fig. 1. (a) Electron-density projection along the a -axis; contours at intervals of l e. \hat{A}^{-2} , with the one-electron line broken. (b) Projection of the structure on (100).

 $B_0=6.0$ Å². The value of R, the usual discrepancy factor, was 33.6% over the observed reflexions.

Refinement proceeded by computing successive (F_o-F_c) syntheses and shifting the atoms to minimize the slopes at the atomic centres, no changes being made in the temperature parameters. After three cycles R had dropped to 15.9% ; an electron-density projection along the a-axis, computed using all the observed $0kl$ reflexions, is shown in Fig. 1.

x-Coordinates

The approximate orientation of the molecule in three-dimensions was then deduced from an examination of the three-dimensional weighted reciprocal lattice, and the x-coordinate of the molecular centre obtained by computing the variation of the calculated structure factors of the $1k0$ and $10l$ reflexions as the molecule was moved along the a -axis. The positional parameters giving best agreement between F_o and F_c for these reflexions were used in computing structure factors for all the $1kl$ reflexions, the origin now being taken midway between the three sets of non-intersecting two-fold screw axes, as given in the *International Tables for X-ray Crystallography* (1952). The discrepancy factor over the observed reflexions was 45.8 %, a discouragingly poor value. Generalized projections were computed using the measured structure amplitudes and calculated phase angles, and new x -coordinates obtained from the relation

tan
$$
2\pi x_j = S_1(y_j, z_j)/C_1(y_j, z_j)
$$
,

where $S_1(y, z)$, $C_1(y, z)$ are the sine and cosine generalized projections. Comparison with the previous x -

Fig. 2. (a) Cosine and (b) sine generalized projections along the *a*-axis; using the lkl data. Contours at intervals of 1 e. A^{-2} , negative contours broken, zero contour omitted.

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Table 1. *Peak heights and curvatures in final cycle*

$_{\rm Atom}$	ρ (e. $\rm{\AA^{-3}}$)		$\partial^2 \rho / \partial x^2$ (e. Å ⁻⁵)		$\partial^2 \rho / \partial y^2$ (e. Å ⁻⁵)		$\partial^2 \rho / \partial z^2$ (e. Å ⁻⁵)	
	Obs.	Calc.	Obs.	Calc.	Obs.	Calc.	Obs.	Calc.
A	4.28	4.27	-13.92	-13.83	-23.68	-23.50	-22.99	-21.54
B	4.49	4.52	-14.84	-14.84	-26.64	-26.18	-25.51	-25.15
$\mathcal C$	4.84	4.80	-14.70	-14.78	-30.33	-29.28	-29.88	-28.53
D	4.78	4.74	-15.35	-15.16	-27.10	-27.69	-27.39	-26.38
\boldsymbol{E}	4.91	4.71	-15.86	-15.09	-28.27	-27.86	-31.08	-28.12
\boldsymbol{F}	4.35	$4 - 25$	-13.71	-13.17	-22.39	-22.44	-26.01	-23.21
G	4.13	4.19	-12.68	$-13-13$	-21.96	-22.12	-24.28	-23.28
A'	3.88	3.98	-11.47	-11.88	-21.65	-21.96	-19.08	-19.33
$_{B^{\prime}}$	4.29	4.36	-13.77	-13.88	-24.05	-24.85	-24.26	-24.83
C'	4.56	4.66	-13.90	-14.32	-28.66	-29.25	-25.42	-25.03
D'	4.76	4.75	-15.60	-14.99	-28.47	-28.25	-29.04	-27.43
E^\prime	4.74	4.78	-16.24	-16.26	-27.30	-27.78	-29.96	-28.53
F^\prime	4.19	4.18	-12.86	-12.53	-21.58	-23.03	-24.44	-22.69
G'	4.09	4.22	-13.28	-13.34	-21.43	-22.37	-22.33	-22.75
P	4.15	$4 - 29$	-12.50	-13.02	-22.21	-23.43	-19.60	-19.11
Ω	4.72	5.05	-13.89	-15.56	-19.64	-20.97	-19.96	-19.86

coordinates showed that there were some quite large changes in fractional coordinate, although these changes did not appear to be quite so drastic when expressed in A units, and the molecular orientation and position were changed only slightly. New coordinates were obtained by applying double shifts, and recalculation of the $1kl$ structure factors revealed that R had dropped to 29.6% , an encouraging improvement. A second cycle reduced R to 22.7% ; the cosine and sine generalized projections at this stage are shown in Fig. 2.

Three-dimensional refinement c

At this stage structure factors were calculated for all the observed reflexions, using the y - and z-coordinates from the *Okl* zone refinement, the x-coordinates from the $1kl$ refinement, and the original temperature parameters $(B_c=5.4, B_o=6.0)$. The R value was 21.3% .

Refinement proceeded by computing observed and calculated differential syntheses, new positional parameters being obtained from the corrections of the F_o synthesis, combined with the back-shift (seriestermination) corrections from the F_c synthesis, and new temperature parameters being chosen to minimize the differences between the observed and calculated peak electron-densities and curvatures. In making the coordinate changes, the shifts applied were 1½ times those obtained from the differential syntheses. Recalculation of the structure factors for all the observed hkl reflexions showed that R had been reduced to 14.9% .

A second set of observed and calculated differential syntheses was computed, and some further small changes in the positional parameters were indicated. The $1\frac{1}{2}$ shift in the previous cycle had apparently been a reasonable value, since some of the atoms had been overshifted and others undershifted. In this second cycle the shifts applied were taken directly from the differential syntheses. The values of the observed and

calculated peak electron-densities and curvatures are listed in Table 1, and these indicate that some further small changes are required in the temperature parameters, especially for the oxygen atom. The final B values adopted are included in Table 2.

Further refinement would require that allowance be made for the contribution of the hydrogen atoms, and for the anisotropic thermal vibrations of the carbon and oxygen atoms. In view, however, of the rather small proportion of reflexions observed (this is reflected in the low peak heights and curvatures) it is doubtful whether such further refinement would be of much value. In any case the required information about the molecular configuration and dimensions can be obtained with fair accuracy at this stage, and so refinement was terminated at this point. The values of the measured structure amplitudes for the observed reflexions are listed in Table 4. To save space the calculated amplitudes and phase angles have not been listed (we actually did not trouble to compute the final

values, but they can of course be readily computed from the positional and temperature parameters of Table 2). Throughout the refinement structure factors were calculated at each stage for many of the unobserved reflexions, and no anomalies were found.

Coordinates and molecular dimensions

The positional and temperature parameters of the carbon and oxygen atoms obtained from the final observed and calculated differential syntheses are listed in Table 2, x , y and z being coordinates referred to the origin of the *International Tables for X-ray Crystallography* (1952), and expressed as fractions of the unit-cell edges.

The best plane through the aromatic carbon atoms, obtained by the least-squares method, is

$$
0.8213X - 0.5039Y - 0.2674Z - 0.7003 = 0,
$$

where X , Y , Z are coordinates expressed in \AA units. The deviations of the atoms from this plane are listed in the final column of Table 2.

Now we must also compute the equation of the plane of the aldehyde group, but since the position of the hydrogen atom is not determined by the X-ray analysis, we will define the required plane as that through the atoms D , P and O . The aldehyde hydrogen atom will lie very close to this plane. The equation is

 $0.5045X - 0.6341Y - 0.5860Z + 0.9951 = 0.$

The angle between the anthracene and aldehyde planes is 27.0° .

The bond lengths and valency angles in the molecule were calculated from the coordinates of Table 2; although there were some quite large variations in the bond lengths $(Fig. 3(a))$, it was not considered (in view of the estimated standard deviations of the atomic positions given below) that the symmetry of the anthracene nucleus differed significantly from *mmm,* and the mean values of the bond lengths and valency angles are shown in Fig. $3(c)$. The standard deviations (see below) suggest that the values of the bond lengths cannot be quoted to better than 0.01 Å, nor the valency angles to better than 1° .

The distance between the oxygen atom and atom F (in a planar model these atoms would approach to within the normal van der Waals distance) is $3.06~\text{\AA},$ and the distance between O and the hydrogen bonded to F (assuming that the hydrogen lies on the anthracene plane, with $C-H=1.08$ Å) is 2.34 Å. (It is possible that H_F is pushed slightly out of the anthracene plane, so that the $O-H_F$ distance may be a little greater than this).

Standard deviations

The standard deviations of the atomic positions were calculated from Cruickshank's (1949) formulae. The r.m.s. values for all the atoms were $\sigma(x) = 0.006$ Å,

Fig. 3. (a) Measured bond lengths. (b) Measured valency angles. (c) Nean bond lengths and valency angles.

 $\sigma(y) = \sigma(z) = 0.005$ Å. Since, however, all the *hkl* structure factors $(h, k, l\neq 0)$ are complex, these values are increased to about 0.01 Å. This is probably a rather optimistic value in view of the large number of unobserved reflexions, and double this value is possibly more realistic. The standard deviations of the mean bond lengths and valency angles are then about 0.010-0.015 Å for the aromatic bonds, 0.015-0.025 Å for the exocyclic bonds, and $1-2^{\circ}$ for the angles.

Discussion

Within the limits of experimental error the anthracene nucleus is completely planar, the maximum deviation of the aromatic carbon atoms from the mean plane being $0.06~\text{\AA}$ and the r.m.s. deviation $0.035~\text{\AA}$. The deviations of the carbon and oxygen atoms of the aldehyde group from this plane $(+0.10 \text{ Å}$ and -0.26 Å respectively) are, however, highly significant.

In a completely planar model there would be unusually close non-bonded distances between the oxygen atom and neighbouring carbon and hydrogen atoms, and the measured molecular dimensions indicate that

the strain due to this overcrowding is relieved by a 27° rotation of the aldehyde group, about the C(aromatic)-C(aldehyde) bond, out of the anthracene plane. In addition there is a displacement of the $C_{\rm al}$ atom out of the anthracene plane in the opposite direction from the oxygen displacement, and a repulsion of the C_{al} and oxygen atoms away from atom F , so that the angles *EDP* and *DPO* are increased from their normal values to 125° and 127° respectively, with the angle *CDP* consequently being reduced to only 111[°]. These deviations from the ideal planar structure increase the distances between the oxygen atom, and atom \overline{F} and its hydrogen atom, to 3.06 Å and 2.34 Å respectively, these distances being almost exactly equal to the sums of the corresponding van der Waals radii (1.7 Å for carbon, 1.4 Å for oxygen and 1.0 -1.2 Å for hydrogen). The distortions of the valency angles from 120° are such that the angle between the planes of the anthracene nucleus and the aldehyde group is reduced to as small a value as is possible without introducing any undue strain due to close intramolecular distances.

The bond lengths in 9-anthraldehyde are compared in Table 3 with those in anthracene (Sinclair, Robertson & Mathieson, 1950; Cruickshank, 1956). Except for bond a, the differences are not significant; for bond a the difference is of the same order of magnitude as that found in 9-nitroanthracene (Trotter, 1959d), but, considering the standard deviations of the present results, the difference is of doubtful significance.

The C(aromatic)–C(aldehyde) bond length (1.46 Å) would appear to be a little shorter than the normal single bond distance in this type of molecule (1.50 Å) , due no doubt to resonance interaction between the aldehyde group and the aromatic rings. It may be noted that the deviations from coplanarity reduce this resonance, but only to 0.79 (cos² 27[°]) of that in a planar model. The $C=O$ bond distance (1.21 Å) is equal to the normal length of this bond in aldehydes.

The fact that the deviations from coplanarity reduce the resonance between the substituent aldehyde group and the anthracene nucleus by only about 20% compared with a planar model explains all the apparent inconsistencies in infrared and ultraviolet spectra and dipole moment data in this type of molecule. The aldehyde group is conjugated to a large extent with the aromatic π -electrons.

It is instructive to compare the deviations from coplanarity with those in the corresponding nitro

Table *4. Observed structure amplitudes*

compounds, 9:10-dinitroanthracene and 9-nitroanthracene. In both of these molecules the environments of the nitro groups are symmetrical, and the strain in planar models is relieved by a 64° rotation of the nitro groups out of the anthracene planes (in crystals of the 9-nitro derivative, crystal forces give an 85° rotation, but the most stable state of the isolated molecule has a 64° rotation (Trotter, 1959 f)). The strain cannot be relieved by distortions of the valency angles or displacement of the nitrogen atoms from the mean planes, for any such distortion or displacement, while reducing the strain at one oxygen atom, would increase the overcrowding at the other oxygen. In 9 anthraldehyde, however, the substituent group is of course not symmetrical, and since the hydrogen atom is smaller than the oxygen the strain in a planar model can be partly relieved by a distortion of the valency angles and a displacement of the aldehydie carbon atom. These quite small variations in valency angles and aldehydic carbon atom position allow the plane of the aldehyde group to make a much smaller angle with the anthracene plane (27°) than the nitro groups in the nitro derivatives (64°) —the corresponding displace-

Fig. 4. Projection of the structure along [100], showing the shorter intermolecular distances.

ments of the oxygen atoms from the aromatic plane are 0.26 Å in the aldehyde and 0.96 Å in the nitro derivatives. The deviations from coplanarity in both types of molecule are such as to make the resonance interaction between the substituent groups and the atomatic rings as large as possible while keeping the non-bonded intramolecular distances equal to the sums of the corresponding van der Waals radii.

Intermolecular distances

All the intermolecular distances correspond to normal van der Waals interactions. The perpendicular distance between the aromatic planes of molecules related by translation a is 3.52 Å. The shorter lateral contacts are illustrated in Fig. 4.

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